

Influences of inner core radius on thermal convection in a rotating spherical shell near the critical Rayleigh number

数値実験に基づく内核半径変化が臨界レイリー数付近での回転球殻対流に与える影響

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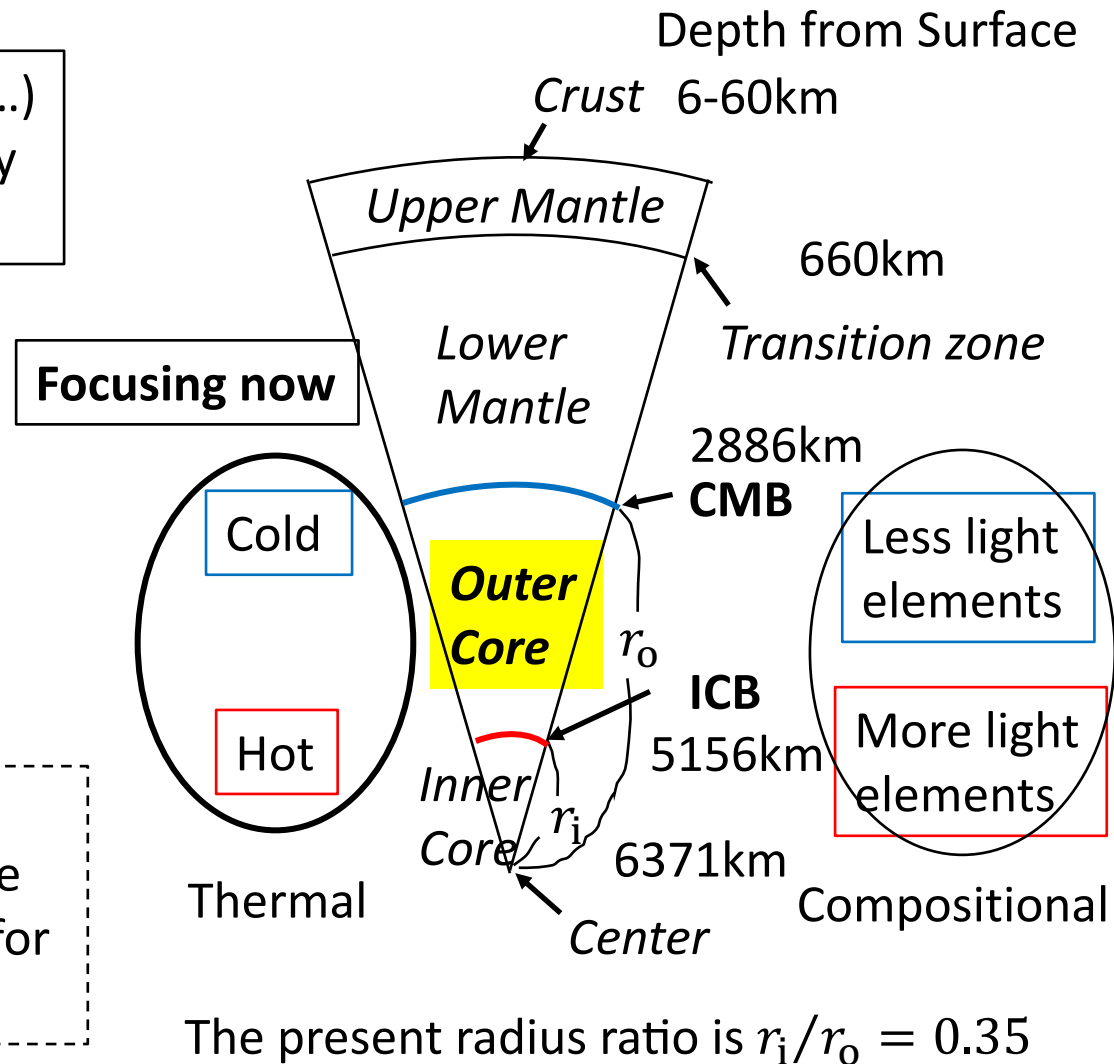
1.1 Focusing on the inner core radii

Convection of fluid Fe (+Si,O,Mg,...)
Kinetic energy → magnetic energy
= “Dynamo process”

Paleomagnetic analysis shows geomagnetic field has been sustained for at least **3.5 billion years**.

+

From **thermochemical evolution** calculation, the inner core could be growing from $r_i/r_o = 0$ to **0.35** for the last billion years.



--> It is important to investigate convection with various inner core radii.

1.2 Governing equations for numerical dynamo

Momentum equation

$$\frac{\partial \mathbf{u}}{\partial t} + (\nabla \times \mathbf{u}) \times \mathbf{u} = -\nabla \left(P + \frac{1}{2} \mathbf{u}^2 \right) - \underbrace{\nabla \times (\nabla \times \mathbf{u})}_{\text{Viscosity}} + \underbrace{\frac{Ra}{Pr} T \frac{\mathbf{r}}{r_0}}_{\text{Buoyancy}} - \underbrace{\frac{2}{E} \mathbf{e}_z \times \mathbf{u}}_{\text{Coriolis}} + \underbrace{\frac{1}{Pm \cdot E} (\nabla \times \mathbf{B}) \times \mathbf{B}}_{\text{Lorentz}}$$

Heat equation $\frac{\partial T}{\partial t} = -(\mathbf{u} \cdot \nabla)T + \frac{1}{Pr} \nabla^2 T$

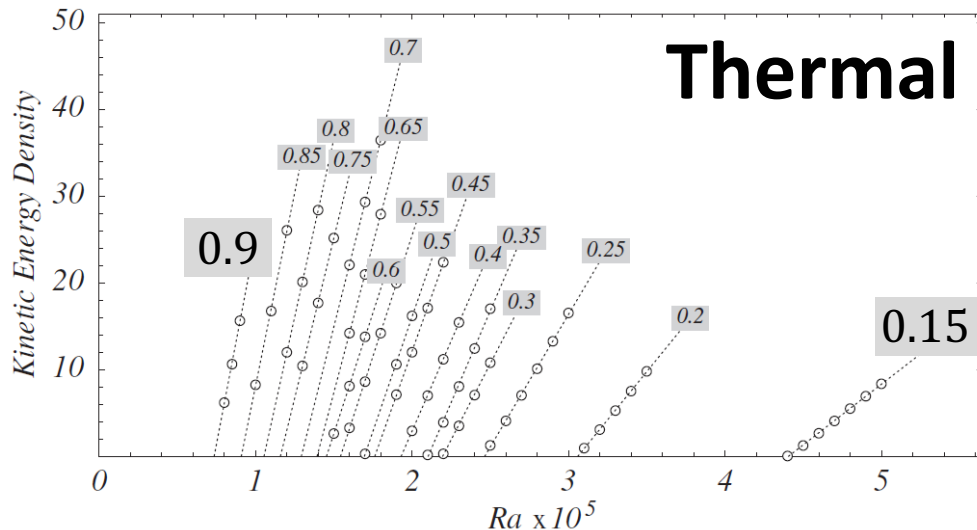
Continuity equation of incompressible fluid $\nabla \cdot \mathbf{u} = 0$

Magnetic induction equation $\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}) - \frac{1}{Pm} \nabla \times (\nabla \times \mathbf{B})$

Gauss's law for magnetic field $\nabla \cdot \mathbf{B} = 0$

Name	Symbol	Definition	Meaning	Estimation
Rayleigh number	Ra	$\alpha_T g_0 (\Delta T) L^3 / \nu \kappa_T$	Buoyancy strength	10^{28}
Ekman number	E	$\nu / \Omega L^2$	Rotation weakness	10^{-15}
Prandtl number	Pr	ν / κ_T	Thermal diffusion weakness	0.1
Magnetic Prandtl number	Pm	ν / η	Magnetic diffusion weakness	10^{-6}

1.3 Previous studies



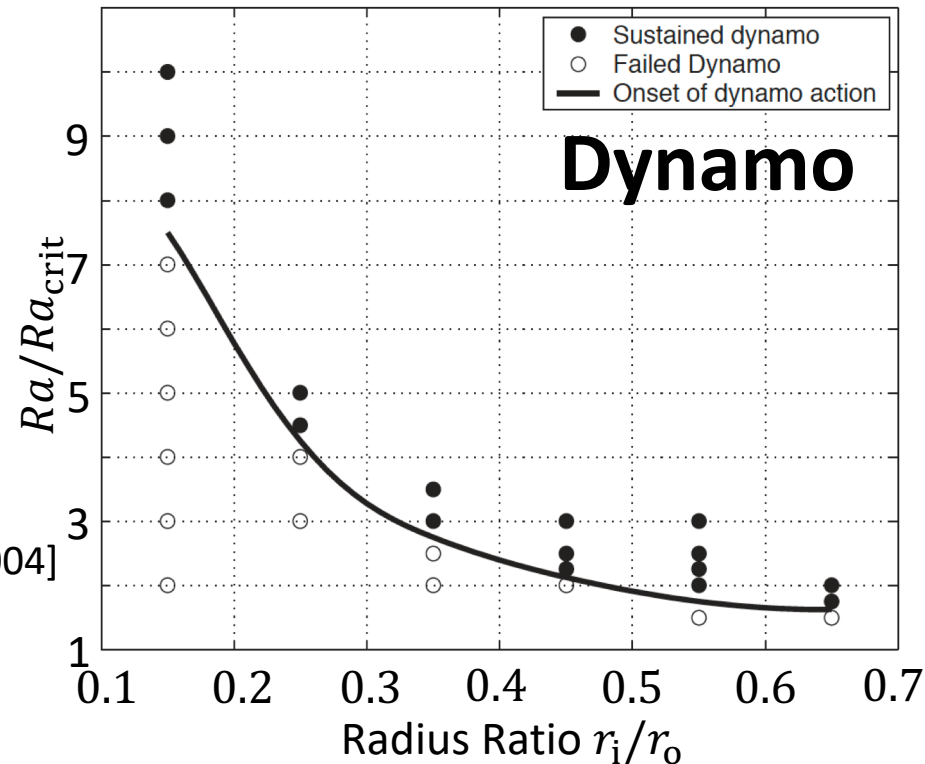
Onset of thermal convection [Al-Shamali et al., 2004]

r_i/r_o : smaller $\rightarrow Ra_{crit}$: bigger

Rayleigh Number (related to buoyancy)

$$Ra = \frac{\alpha_T g_0 (\Delta T) L^3}{\nu \kappa_T}$$

*The critical Rayleigh number Ra_{crit} represents onset of thermal convection.



Dynamo onset [Heimpel et al., 2005]

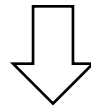
r_i/r_o : smaller
 $\rightarrow Ra$ needed to sustain dynamo: larger

➤ Influence of inner core radius on convection is not fully understood.

1.4 Purpose of this study

For fully understanding of the convection in the outer core at **the past Earth**, we investigate properties of dynamo action with the smaller inner core using numerical dynamo open code Calypso [Matsui et al., 2014].

- Calculation of **kinetic/magnetic energy** density
- Calculation of **length scale** of flow



- **Maximum growth mode** in magneto-convection model
- Influence of magnetic field on convection near dynamo onset

2.1 Code and initial/boundary condition

We used numerical dynamo open code Calypso [Matsui et al., 2014].

- radial derivatives ... second order finite difference
- spectral method
 - solenoidal vector field ... poloidal and toroidal
- time stepping
 - the linear diffusive terms ... the Crank-Nicolson
 - the Coriolis force and the nonlinear terms ... second order Adams-Bashforth

Initial condition [cf. Benchmark case 1 proposed by Christensen et al. (2001)]

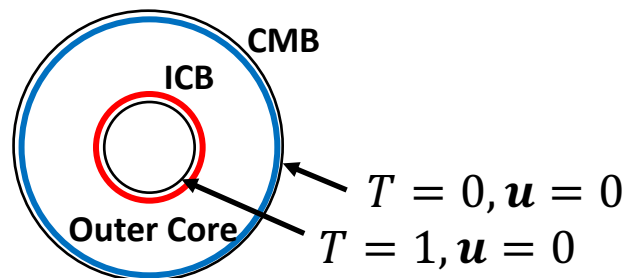
Given **various modes** of temperature perturbation by

$$T(r, \theta, \phi) = \sum_{l=1}^{l_{\max}} T_l^l(r) Y_l^l(\theta, \phi)$$

Given **an axial dipole** as a seed magnetic field;

$$\mathbf{B} = \nabla \times \nabla \times (B_{S1}^0(r) Y_1^0(\theta, \phi) \hat{r}) + \nabla \times (B_{T2}^0(r) Y_2^0(\theta, \phi) \hat{r})$$

Boundary condition



Mantle and the inner core
...Electrically insulated

2.2 Parameters setting

	Case 1	Case 2	Case 3
r_i/r_o	0.15	0.25	0.35
$Ra_{crit}[\times 10^5]$	1.09	0.72	0.56
$Ra[\times 10^5]$	8.7~17	1.4~7.0	0.84~4.0
Ra/Ra_{crit}	7.0~15.6	1.9~9.7	1.5~7.1

<-- Thermal simulation

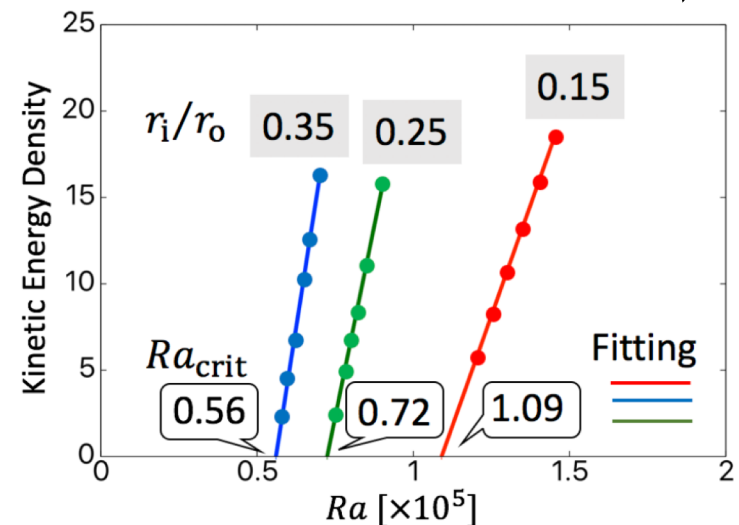
*The critical Rayleigh number Ra_{crit} represents onset of thermal convection.

$$E = 1 \times 10^{-3}, Pr = 1, Pm = 5$$

Kinetic/Magnetic energy density

$$E_{kin} = \frac{1}{2V_S} \int_{V_S} \mathbf{u}^2 dV$$

$$E_{mag} = \frac{1}{2V_S E Pm} \int_{V_S} \mathbf{B}^2 dV$$



3.1 Kinetic/magnetic energy density

Time average of E_{kin} and E_{mag} for $t/\tau_\eta = 1.5$ to 2.0

* $\tau_\eta (= L^2/\eta)$: magnetic diffusion time

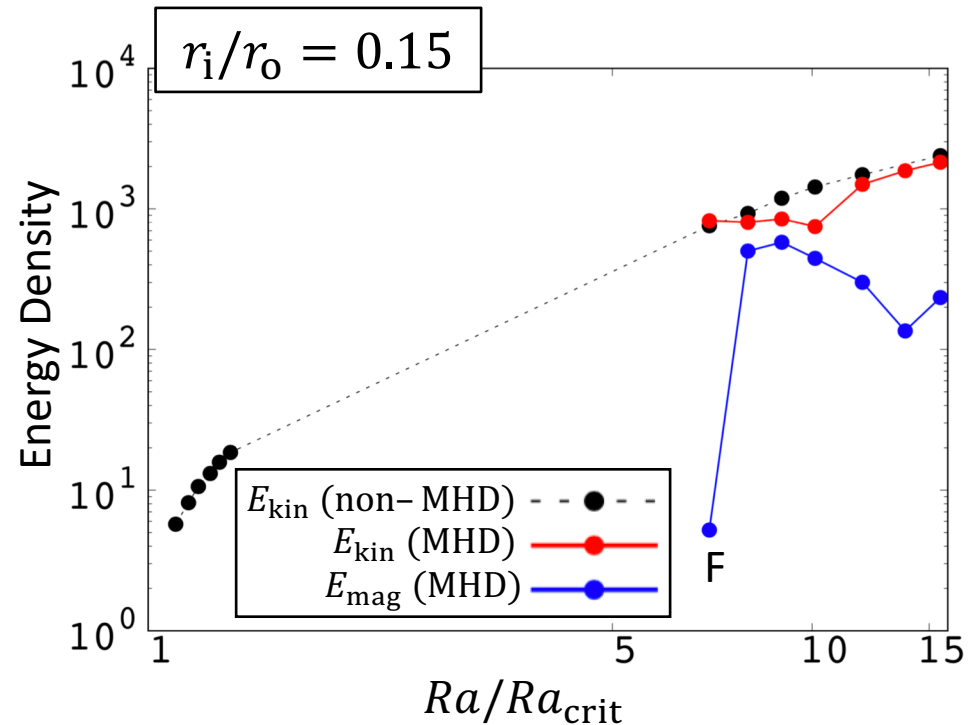
➤ Common point

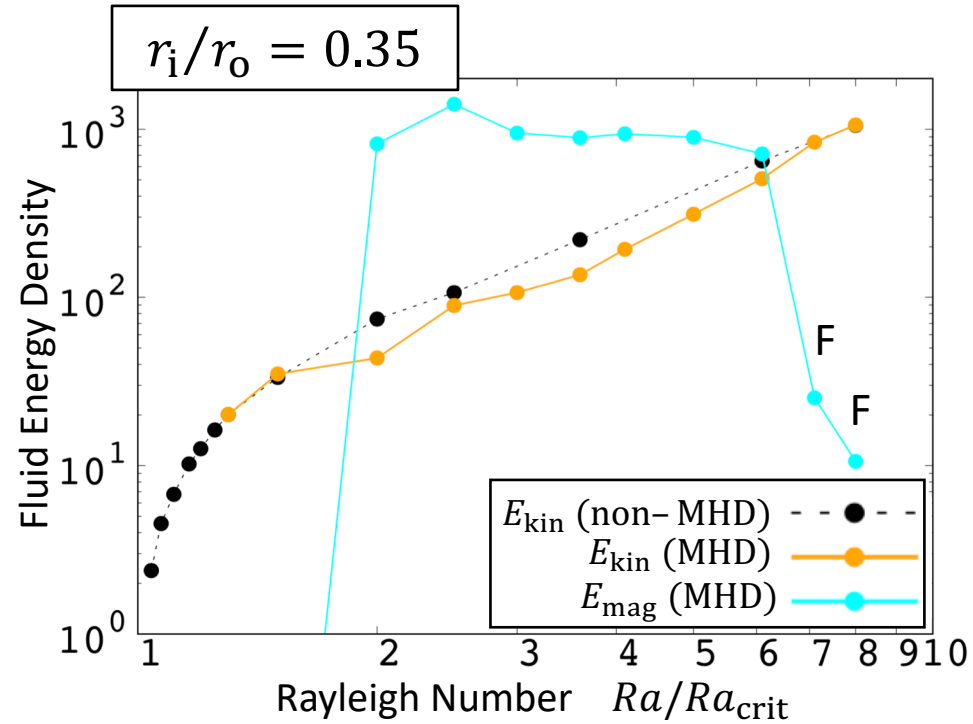
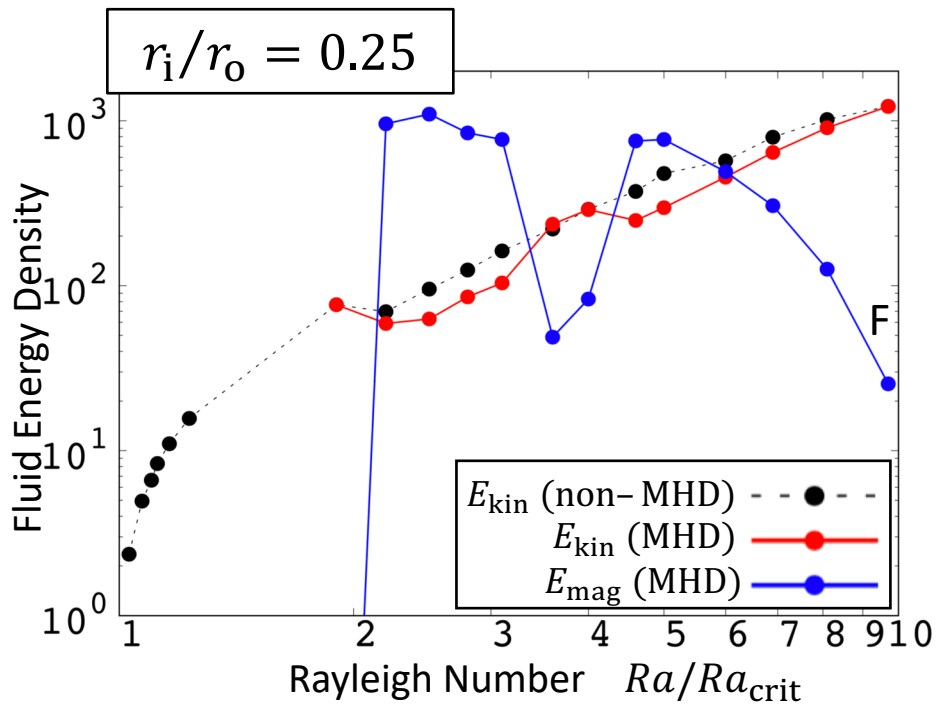
E_{kin} (MHD) $<$ E_{kin} (non-MHD) under large magnetic field.

➤ Deferent points

- In $r_i/r_o = 0.15$, $E_{\text{mag}} < E_{\text{kin}}$
- In $r_i/r_o = 0.25$, $E_{\text{mag}} > E_{\text{kin}}$ or $E_{\text{mag}} < E_{\text{kin}}$
- In $r_i/r_o = 0.35$, $E_{\text{mag}} > E_{\text{kin}}$

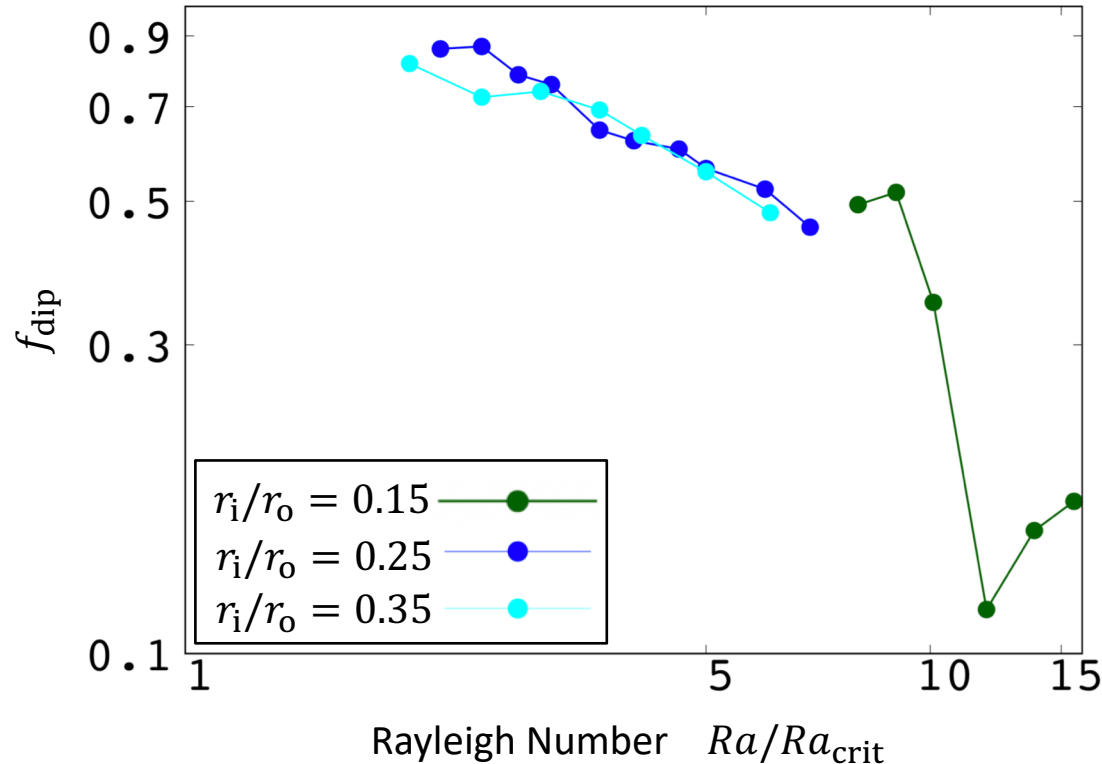
Lorentz force is strong for $E_{\text{mag}} > E_{\text{kin}}$.





- E_{mag} drops at $Ra/Ra_{crit} = 3.1$ and 3.6 in $r_i/r_o = 0.25$.
- E_{mag} is comparable in $r_i/r_o = 0.35$.

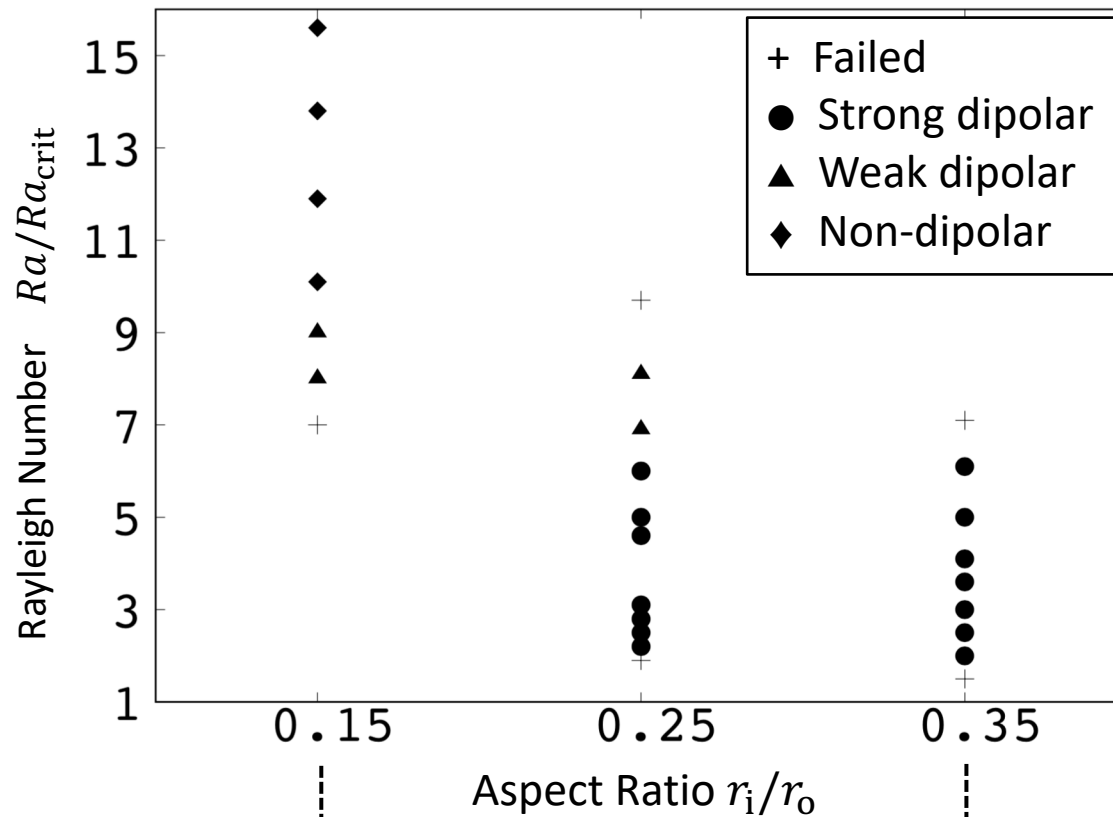
3.2 Dipolarity



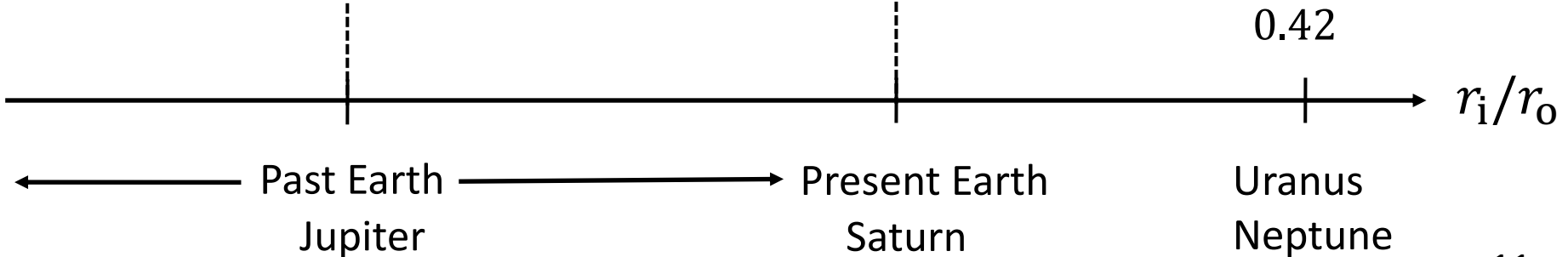
Dipolarity is an index which represents strength of an axial dipole component.

$$f_{dip} = \left(\frac{E_{mag}^{(l=1,m=0)}(r = r_o)}{\sum_{l=1}^{l_{max}} \sum_{m=0}^l E_{mag}^{(l,m)}(r = r_o)} \right)^{1/2} \quad \left[\begin{array}{l} l: \text{Spherical harmonic degree} \\ m: \text{Spherical harmonic order} \end{array} \right]$$

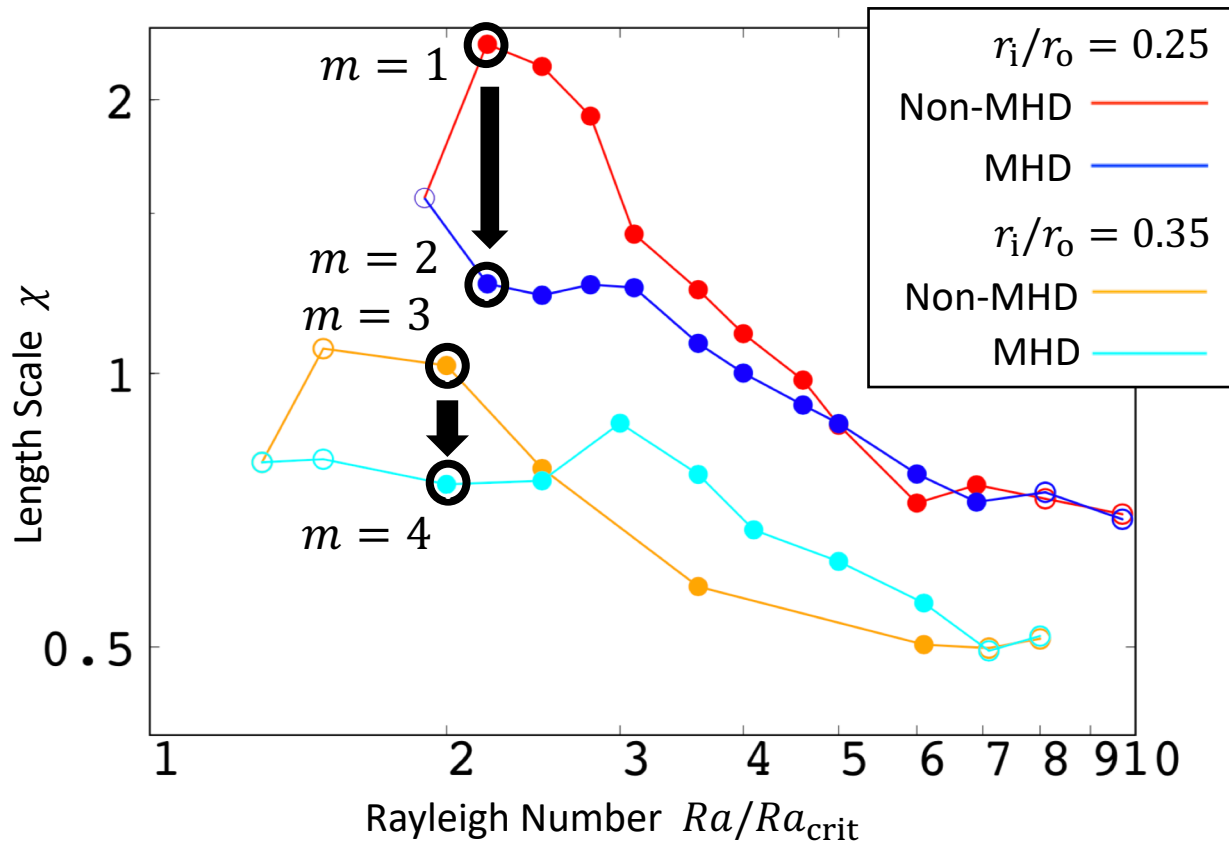
3.3 Dynamo regime in various aspect ratios



- When Ra becomes large, an axial dipolar component becomes small.
- In the smaller inner core size cases, non-dipolar components become large.



4.1 Length scale of flow in azimuthal direction



$$\chi = \frac{\overline{\pi < \mathbf{u}^2 >}}{\sum m < \mathbf{u}_m^2 >} L = \left\{ \pi / \left(\sum m \frac{\overline{< \mathbf{u}_m^2 >}}{< \mathbf{u}^2 >} \right) \right\} L \quad [\text{cf. King and Buffett, 2013}]$$

(m : Spherical harmonic order, \mathbf{u} : velocity, L : the outer core thickness)

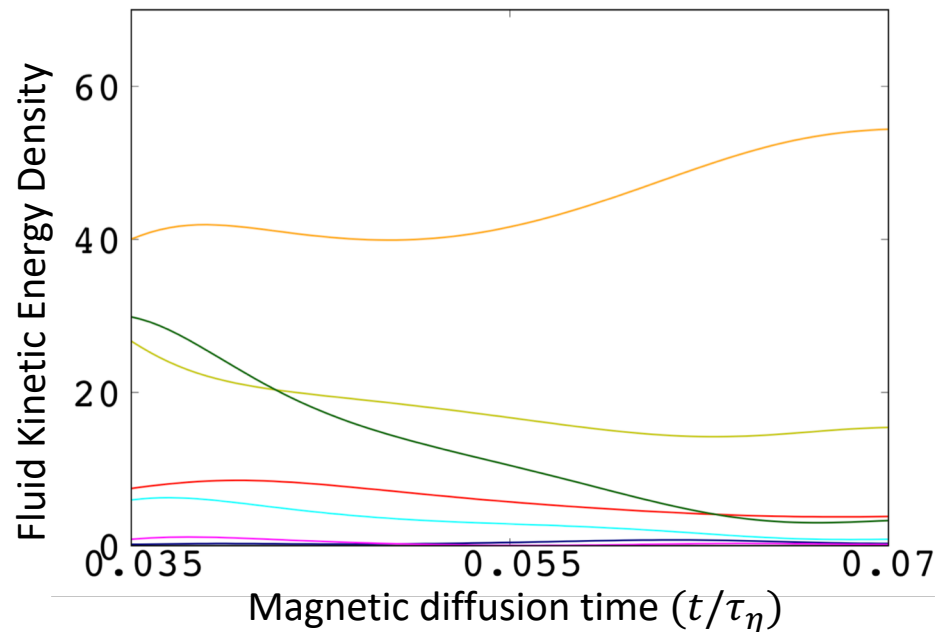
4.2 Growth in different modes

Simulation using linear terms (under background magnetic field and $\partial \mathbf{B} / \partial t = 0$)

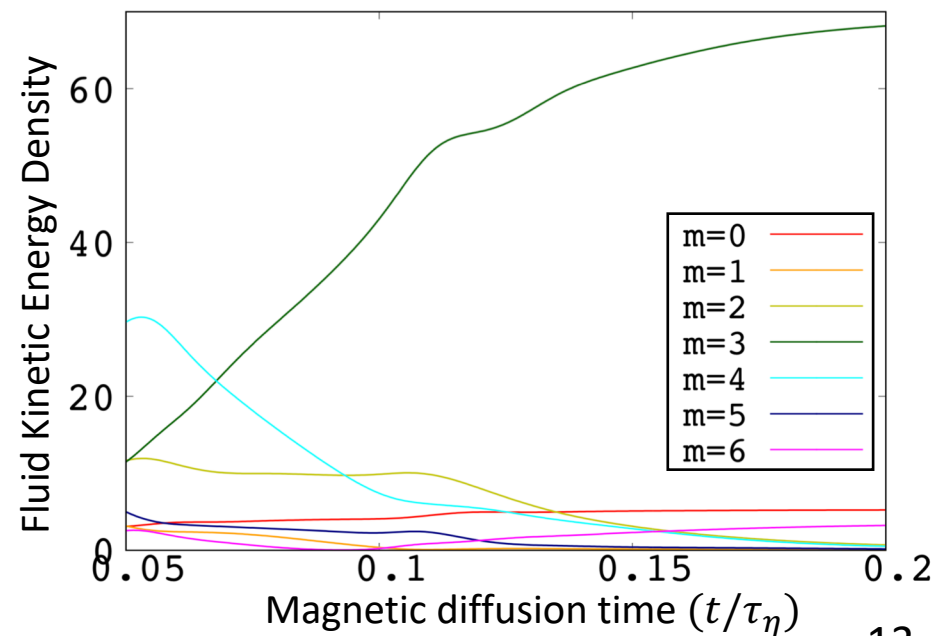
$$\frac{\partial \mathbf{u}}{\partial t} + \cancel{(\nabla \times \mathbf{u}) \times \mathbf{u}} = -\nabla \left(P + \cancel{\frac{1}{2} \mathbf{u}^2} \right) - \nabla \times (\nabla \times \mathbf{u}) + \frac{Ra}{Pr} T \frac{\mathbf{r}}{r_0} - \frac{2}{E} \mathbf{e}_z \times \mathbf{u} + \frac{1}{Pm \cdot E} (\nabla \times \mathbf{B}) \times \mathbf{B}$$

$$\frac{\partial T}{\partial t} = -\cancel{(\mathbf{u} \cdot \nabla) T} + \frac{1}{Pr} \nabla^2 T \quad \nabla \cdot \mathbf{u} = 0$$

$$r_i/r_0 = 0.25, Ra/Ra_{\text{crit}} = 2.2$$



$$r_i/r_0 = 0.35, Ra/Ra_{\text{crit}} = 2.0$$



4.3 Maximum growth mode

Maximum growth mode near dynamo onset is

- $m = 1$ (no B) \rightarrow $m = 1$ (with stable B) $\rightarrow m = 2$ (with B) in $r_i/r_o = 0.25$
- $m = 3$ (no B) \rightarrow $m = 3$ (with stable B) $\rightarrow m = 4$ (with B) in $r_i/r_o = 0.35$

Estimated growth rate p from the time history of E_{kin} which is set as $\exp(pt)$

m	0	1	2	3	4	5	6
$p (r_i/r_o = 0.25)$	-32.41	68.30	-84.09	-93.97	-40.93	-90.08	-448.8
$p (r_i/r_o = 0.35)$	7.62	-2.96	-24.91	27.03	-22.48	-27.47	13.47

5 Summary and future work

- In the **smaller inner core** size cases, **non-dipolar** components become large.
- **Length scale** of flow in MHD cases is smaller than that in non-MHD cases with small Ra .
- **Maximum growth mode** in MHD cases is larger than that in magneto-convection model.

✓ Future works

- More simulations in some Ra in $r_i/r_o = 0.25, 0.35$
→ Investigation of maximum growth mode dependency on Ra and radius ratio and initial field
- Derivation of the maximum growth mode with/without the inner core from linear stability analysis... scale of Ra_{crit} on Ekman number and aspect ratio
[Bissopp, 1958; Chandrasekhar, 1961; Roberts, 1968; Busse, 1970; Jones et al., 2000]